

Detrimental effect of self-phase modulation on the performance of Brillouin distributed fibre sensors

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We show that the spectral broadening of the pump pulse through self-phase modulation in a time-domain distributed Brillouin sensor has a considerable detrimental effect in the measurement, especially in the case of long distances and high resolution pulses. Using 30 ns pump pulses with peak power of ~350 mW, self-phase modulation leads to a doubling of the effective gain linewidth after some 20 km, which is equivalent to a contrast loss of 2 dB in the measurement. The impact is higher for shorter pulses (higher resolution). The theoretical modeling is fully confirmed by experimental results. © 2010 Optical Society of America
OCIS codes: 290.5900, 190.2640, 060.2370, 190.4370, 190.3270.

For at least two decades Brillouin fibre sensors have attracted a great interest in the fibre sensing community for their temperature and strain monitoring capability [1,2]. In time-domain distributed Brillouin sensors pulses are used to interrogate the local interaction in the fibre. The accuracy on the measurand is scaled by the spectral spreading of the effective gain, which at its turn is given by the convolution between the pulse spectrum and the natural Brillouin gain spectrum (BGS). According to standard time-bandwidth relations, the Gaussian pulse is presumably the best candidate for this interrogation when compared to other profiles (rectangular, triangular, etc.). However, we show here that this is not the best choice when addressing long-ranges, because this pulse shape leads to a significant spectral broadening of the BGS along the fibre. An observed broadening of the BGS was suspected to be caused by self-phase modulation (SPM) in an early work by Lecoecue *et al* [3] and by Izumita *et al* [4] in coherent OTDR system. SPM leads to small phase chirps during intensity transitions in the pump pulse (leading and trailing edges) that eventually become important in long fibres. The frequency broadening associated with this phase modulation leads to a reduced peak gain and uncertainties in the determination of the Brillouin shift ν_B , but leaves the temporal intensity distribution of the pump pulse unchanged and hence the spatial resolution is preserved. While the former [3-5] works showed a correct intuition addressing qualitatively the issue, no theoretical model was given or the model was incomplete [5]. We present here a quantitative model of the detrimental impact of SPM supported by a clear experimental demonstration. Two optical pulses with different temporal profiles were judiciously chosen (rectangular and Gaussian) showing the same full width at half maximum (FWHM) and carrying the same energy, to evenly study and compare the SPM impact on their spectrum. Then we clearly experimentally demonstrate the spectral broadening of the BGS due to SPM in a Brillouin distributed sensor in various conditions in terms of pump pulse temporal

profiles, power and width. The results are compared with a theoretical model showing good agreement.

SPM is a consequence of the nonlinear Kerr effect in the fibre that results in an intensity-dependent refractive index. It leads an optical pulse to modulate its own optical phase according to its intensity profile. Since the pulse creates a time-dependent refractive index variation, higher for larger intensities, it leads to a time-dependent nonlinear phase shift ϕ_{NL} that translates into variations of the instantaneous frequency (frequency chirp) along the pulse. Let us consider the propagation of an optical pulse through a fibre characterized by a nonlinear refractive index $n(t)=n_0+n_2 I(t)$. After propagation over a distance z , the optical pulse accumulates an additional nonlinear phase [6] due to the intensity-dependent part $\phi_{NL}=n_2 (\omega/c_0) z I(t)$ of the refractive index which creates a variation in its instantaneous frequency.

$$\Delta\omega(t)=\frac{d\phi_{NL}(z,t)}{dt}=-n_2 \frac{\omega}{c_0} z \frac{dI(t)}{dt} \quad (1)$$

Eq. 1 clearly shows that SPM will have more impact on fast signals constantly showing temporal transitions, with no interval of constant intensity. For our study we have compared two different temporal profiles (rectangular and Gaussian) showing the same width (for equal spatial resolutions) and the same energy (to make gains equal). The rectangular pulse offers constant intensity intervals only while the Gaussian pulse varies constantly in intensity but shows the best time-bandwidth product (and thus the best time-frequency resolution). Since the intensity is constant at any time for the rectangular profile (assuming instantaneous transitions), $\Delta\omega(t)=0$ and no pulse spectral broadening should therefore be observed. The situation will be drastically different in the Gaussian case showing continuous time transitions. Since it can be entirely described analytically, let consider a pulse with a Gaussian envelope, $E(0,t)=A \exp(-t^2/\tau^2)$, so,

$I(t)=|A \exp(-t^2/\tau^2)|^2=I_0 \exp(-2t^2/\tau^2)$ with a $1/e$ width τ . The instantaneous frequency chirp imposed on the Gaussian pulse by SPM using Eq.1 can be find in Eq.2 as:

$$\Delta\omega(t)=4n_2 \frac{\omega}{c_0} z I_0 \frac{t^2}{\tau^2} \exp(-2t^2/\tau^2)=4\gamma z P \frac{t^2}{\tau^2} \exp(-2t^2/\tau^2) \quad (2)$$

where c_0 is the vacuum light velocity, ω is the central angular frequency of the optical pulse, I_0 and P are, respectively, the peak intensity and peak power of the optical pulse, A_{eff} is the effective area of the fibre mode and $\gamma=n_2\omega/c_0A_{eff}$ is the fibre nonlinear coefficient. In a long fibre segment the distance z must be replaced by the effective nonlinear distance $z_{eff} = [1 - \exp(-\alpha L)] / \alpha$ where L is the physical length and α in the linear loss of the fibre. The peak excursion of the instantaneous frequency is $\Delta\omega_{max} = 2\gamma P z_{eff} / \sqrt{e} \tau = 1.43 \gamma P z_{eff} / T$ offset by $\tau_{max} = T / \sqrt{8 \ln 2}$ from the pulse centre and, as scale factor, gives a good estimate of the real spectral broadening. For typical values in conventional long-range BOTDA systems ($\gamma=1 \text{ W}^{-1} \cdot \text{km}^{-1}$, $P=100 \text{ mW}$, $z_{eff}=20 \text{ km}$ and $T=10 \text{ ns}$), this quantity amounts to $2\pi \cdot 29 \text{ MHz}$, which is similar to the Brillouin natural linewidth.

A more precise evaluation of the spectral broadening can be obtained by computing the pulse spectrum evolution along the fibre. The actual frequency spectrum of the Gaussian pulse after experiencing SPM in the fibre is given by the Fourier transform of the output pulse amplitude according to the Eq.3:

$$g_{SPM}(\Delta\nu)=FT \left\{ \underbrace{A \exp(-t^2/\tau^2)}_{\text{Gaussian amplitude term}} \underbrace{\exp[i\gamma L_{eff} P \exp(-2t^2/\tau^2)]}_{\text{SPM term}} \right\} \quad (3)$$

where, for convenience, the time origin is considered always placed at the pulse centre. The effective Brillouin gain linewidth is given by the convolution between the real pump source spectrum after SPM, $g_{SPM}(\Delta\nu)$, and the natural BGS. It must be noted that all along this derivation we have neglected the effect of chromatic dispersion. This is due to the comparatively long pulse lengths used in BOTDA systems (tens of ns). It can be easily shown that in these conditions the dispersion length far exceeds the nonlinear length and thus dispersion effects can be neglected [6]. In relation to this, modulation instability (which results from SPM phase-matched by anomalous dispersion) is known to impose a limit in the peak power of the pump pulses used in Brillouin sensors due to nonlinear pump depletion [7]. The observed effect here is radically different from the one observed in [7] since it depends not only on the pulse power but also on the pulse shape and duration, and causes a smooth broadening of the BGS which was not observed in [7].

The experimental configuration, shown in Fig. 1, is based on a minor adaptation of the high performance pump and probe set-up described in [8]: A 1552 nm DFB laser is used as the light source. Pump and signal waves are derived from the same optical source by means of optical modulation and inserted in the fibre through opposite ends. The pump pulse is shaped by an intensity modulator driven by an arbitrary waveform generator. The pulses produced have 30 ns FWHM optical pulse duration corresponding to 3 m spatial resolution. This duration was carefully chosen for a better demonstration

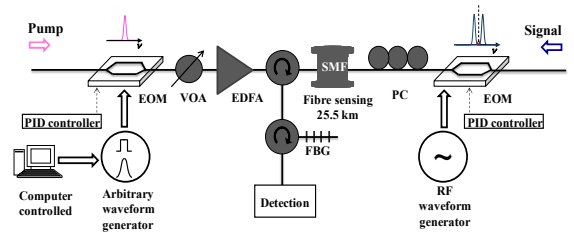
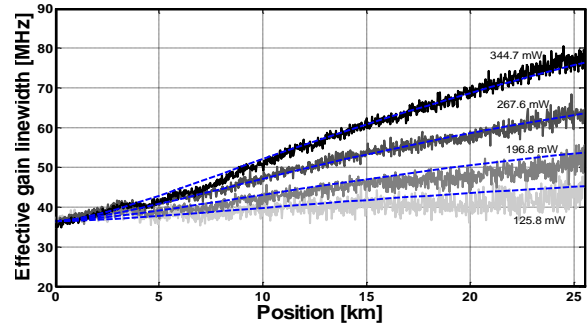
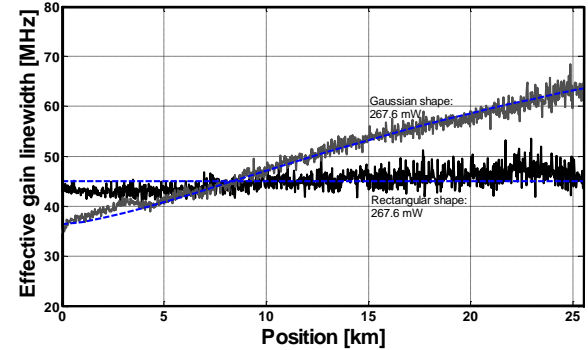


Fig. 2: Diagram of the setup: EDFA: erbium-doped fibre amplifier; EOM: electro-optic modulator; FBG: fibre Bragg grating (0.1 nm spectral width); PC: polarization controller; VOA variable optical attenuator; PID: proportional-integrator circuit.

as a trade-off between temporal transitions fast enough to show the SPM effect and long enough to keep the original spectrum sufficiently narrow (to avoid an excessive broadening of the effective BGS). The output light is controlled by a polarization controller to best align pump and signal polarizations in order to maximize the SBS interaction, and then boosted by a high power EDFA. The BGS is determined by scanning the microwave generator frequency around the Brillouin shift ν_B . The net Brillouin gain is then measured using a detection stage comprising a 125 MHz bandwidth photoreceiver and a digital oscilloscope synchronously triggered by the arbitrary waveform generator. The experimental study of the SPM impact on Brillouin sensors has been carried out through a 25.5 km standard SMF fibre using the above described BOTDA system. The distributed nature of the measurements makes possible to retrieve the effective BGS at any position along the fibre, and thus to observe the gradual impact of SPM. Depletion is



(a)



(b)

Fig. 1: Experimental demonstration (solid lines) and theoretical analysis (dotted lines) of the gain spectrum broadening due to SPM along 25.5 km SMF fibre (a) for different peak powers of a 30 ns FWHM Gaussian pulse (b) comparing 30 ns FWHM Gaussian and rectangular pulses of identical energy and same pump peak power.

avoided by using sufficiently low probe power levels (in the microwatt level) to guarantee depletion values below 1%.

Fig. 2(a) shows the effective gain linewidth (solid lines) along the fibre using 30 ns FWHM Gaussian pulses for increasing pump peak power from 125.8 mW to 344.7 mW. The graph shows clearly a quasi-linear increase in the effective gain linewidth caused by the pump spectral broadening due to SPM along the fibre. The same figure represents the calculation of the effective gain linewidth (FWHM) as a function of the distance (dotted lines) showing a good match between the experimental result and the theoretical model with $\gamma=0.7 \text{ W}^{-1}\text{km}^{-1}$. This value of γ is small in comparison with the values reported usually in conventional SMF ($1.1 \text{ W}^{-1}\text{km}^{-1}$), the discrepancy probably coming from an imperfect power calibration in our setup. Fig. 2(b) represents the evolution of the effective gain linewidth along the fibre using 30 ns FWHM rectangular and Gaussian pulses of identical energy (267.6 mW Gaussian peak power). It can be observed that the gain spectral width is doubled after 25.5 km for the Gaussian pulse, while it remains nearly unchanged for the rectangular pulse as predicted by the model (dotted lines). The residual broadening of the spectrum in the square pulse is probably due to a non-ideal shape in the rectangular pulse used (the rise and fall edges are not ideally sharp). This incidentally confirms a negligible pump depletion that would instead broaden the effective linewidth of the gain spectrum near the fibre end, independently of the pulse shape. The benefit of a better time-bandwidth product using the Gaussian pulse turns out to be entirely cancelled by SPM after less than 10 km. From the Fig. 2(a), and considering the

extreme case of 344 mW peak power, the contrast loss can be estimated at the output of the fibre (25.5 km) to be 2 dB, which is equivalent to 10 km distance penalty in standard conditions.

Fig. 3(a) represents comparative measurements of the gain spectral width measured for a Gaussian pulse at the fibre input (dark gray diamonds) and output (light gray squares), respectively, for different peak powers, demonstrating the clear signature of SPM by the absence of effect at short distance and through the linear dependence on power at long distances. As shown in Fig. 3(b) the effective Brillouin gain linewidth varies in inverse proportion to the optical pulse duration T at both the fibre input and output, but with a different slope. An excess broadening is present at the output as a result of the pump spectral spreading. This excess broadening becomes larger for shorter pulses, showing also a $1/T$ slope, in good agreement with the theory above.

In conclusion, we have demonstrated that SPM-induced spectral broadening can have a significant effect on the effective gain linewidth measured by Brillouin sensors. Modeling and experiments have undoubtedly demonstrated that the effective gain linewidth can easily experience a two-fold increase in standard conditions when the pulse intensity profile is Gaussian. At power and spatial resolutions used in standard configurations, spectral broadening can be observed typically from a 5 km distance. Practically the problem can be circumvented by using a clean rectangular pulse with very sharp rising and falling edges. The benefit of a narrower spectrum for an identical pulse width brought by the Gaussian pulse is rapidly cancelled after about 10 km for 267.7 mW pump peak power.

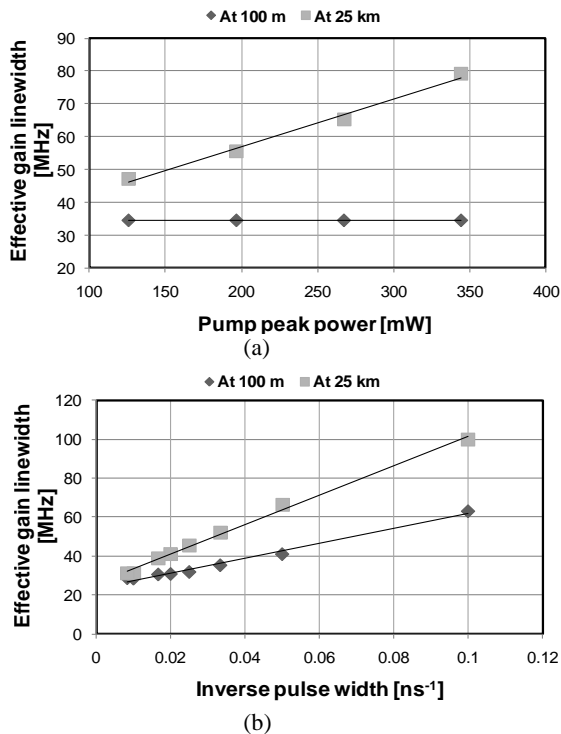


Fig. 3: Brillouin gain spectral width measured close to the fibre input (100 m) and at the fibre output (25.5 km) using a 30 ns Gaussian pulse (a) Linear dependence on the pump peak power, showing that the effect of SPM is observed only at the distant end, (b) Inverse dependence of the gain linewidth on the pulse width, showing the excess broadening due to SPM at the far end (pump power is 196.8mW).

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