High precision and tunable multi-wavelength fiber source based on cascaded four-wave mixing enhanced by Raman

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ABSTRACT

We present a multi-wavelength fiber source based on cascaded of four-wave mixing in two semiconductor optical amplifiers followed by further four-wave mixing in an optical fiber enhanced by Raman amplification. The multi-wavelength source is generated by two initial frequencies detuned 200 GHz and referenced in the absorption lines of the acetylene $^{12}C_2H_2$, which sweep in frequency keeping the detuning of the lasers constant. With this configuration, we have achieved a high resolution source with a spectrum of 36 channels centered with adjustable peaks separation. The source can be employed to interrogate a fiber Bragg grating sensors network and in gas spectroscopy applications.

Keywords: multi-wavelength fiber laser, standard frequencies, four-wave mixing, semiconductor optical amplifier, Raman amplification

1. INTRODUCTION

Multi-wavelength laser sources present a great interest in optical sensing, optical spectroscopy, microwave signal processing, and high capacity WDM optical fiber communications. Depending on the application, the requirements for multi-wavelength sources include a large number of channels over a broad bandwidth, high output power uniformity distributed over the channels, precise position of the ITU frequency grid, high stability of their wavelength emissions or high tunability. Multi-wavelength sources are usually designed using several gain mechanisms such as erbium-doped fiber, semiconductor optical amplifier and stimulated Raman scattering, and channel spacing mechanics as diffraction gratings, Bragg gratings, Lyot filters, Sagnac filters, intracavity etalons or Fabry-Perot micro-etalons [1-7].

An interesting application of the multi-wavelength fiber source is in optical fiber sensors. Nowadays, optical fiber sensors are employed in many applications including strain measurements in civil infrastructure, oil pipeline monitoring, medical applications or chemical sensing. In many of these applications a set of Bragg gratings that forms a network much be interrogated. In many of these applications high resolution, high precision and tunability is required [8].

In this paper, we present a multi-wavelength source for interrogating a fiber Bragg grating (FBG) network using two frequencies obtained by two tunable lasers, separated 200 GHz. Both lasers sweep over the wavelength range of visible absorptions in an acetylene cell, so that there is a continuous absolute wavelength monitoring of the whole source. In order to obtain the multi-wavelength laser source the two lasers are combined to produce degenerate four-wave mixing (FWM) in a cascade of two semiconductor optical amplifiers (SOAs) and enhanced by further FWM and Raman amplification in a dispersion-shifted optical fiber. The input lasers are centered close to the zero dispersion wavelength of the optical fiber used in the experiment in other to increase the FWM efficiency, leading to the generation of up to 36 referenced lines.

2. EXPERIMENTAL SETUP AND RESULTS

Fig 1 shows the experimental setup used to generate the proposed multi-wavelength fiber laser. It consists of two tunable lasers, two cascade SOAs with a saturation power of 13 and 15 dBm respectively, two isolators, two Raman fiber laser at

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wavelength 1455 nm and 1427 nm respectively, a 5 Km dispersion-shifted fiber, an absorption acetylene cell (traceable to NIST [9]) and two optical detectors.

The system works as follows: the output of both lasers are divided in two branches. One is used to reference the laser line to the selected frequency by means of the acetylene cell, and the others are coupled for the multi-wavelength source generation. Initially, TL1 emits at the frequency of the absorption line P23, f_1 =194.743 THz (λ =1539.46 nm) with output power P_1 = 5 mW. TL2 signal is detuned 200 GHz from TL1, so that f_2 = 194.943 THz (λ =1537.85 nm) with P_2 = 6 mW. This frequency spacing can be changed depending on the requirements of the FBG network system to be interrogated. TL1 and TL2 sweep simultaneously keeping the constant detuning, and the sweep is stopped when f_1 arrives at line P27 of the acetylene absorption spectrum (f=194.386 THz, λ =1542.25 nm). Hence there is absolute references for TL1 in the beginning and end of the sweep, and for TL2 these appear along the sweep. Since the spectrum of acetylene is well-known, the position of these references is clearly determined along the sweep and is used as a feedback to correct any errors.

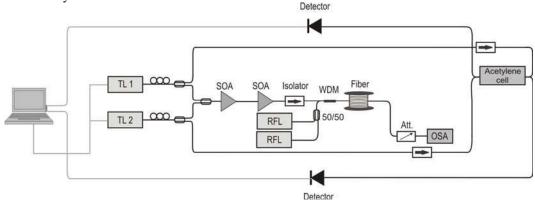


Fig 1.- Experimental setup

The signals emitted from the acetylene cell are detected by two detectors and controlled by a computer. When TL1 or TL2 emissions pass through an absorption line, the computer uses the signal detected by the detectors to correct possible frequency drifts by acting on the sweep signal of f_1 and f_2 . By knowing the absorption spectrum of the acetylene [9] and assuming that the sweep is linear, the frequencies emitted from TL1 and TL2 are known in the whole process.

The signals from TL1 and TL2 are polarization-controlled to maximize the FWM efficiency and are combined at the input of the first SOA. We use two SOAs (Kamelian OPA-20-N-C-FA and Covega BOA-1004) that have longer active region than conventional ones used for amplification, and hence exhibit a larger nonlinear behavior. The two SOAs are cascaded in order to increase the output power and so as to enhance the FWM generation. Fig 2 shows the spectra obtained at the output of the first and second SOA, showing a consistent increase in the number of FWM products generated using two SOAs (nine frequencies) instead of only one SOA (six frequencies).

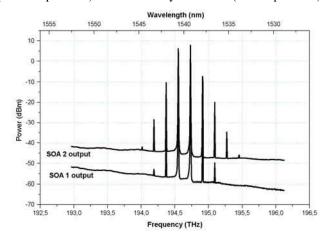


Fig 2.- FWM generated in a high non-linear SOA and enhanced in a high output saturation power SOA

The output of the second SOA is introduced in a 5 km long dispersion-shifted fiber with a zero-dispersion wavelength at 1540.9 nm and a dispersion slope of 0.056 ps•nm⁻²•km⁻¹. The TL1 and TL2 laser frequencies have been selected close to the zero dispersion of the fiber in order to improve the FWM generation in the fiber, and, due to its reduced effective area, the fiber also presents a relatively high Raman gain (3.3 dB·W⁻¹·km⁻¹). The Raman gain is obtained in our experiment with a Raman laser centered a 1455 nm. The benefits of introducing Raman gain in this experiment are double: first, since the input waves are amplified traveling in the fiber, the efficiency of the FWM process is largely enhanced; second, the phase matching conditions of FWM are less restrictive, as it was shown in [10]. As a consequence of this, the intensity of the generated FWM products remains basically constant over the full sweep range, which may be enough to scan the spectrum of a whole fiber Bragg grating

Fig 3 shows the output spectra obtained at the output of the fiber for different Raman pump powers. The number of visible lines increased with the Raman pump power. The optimum results are obtained for a pump power of 1 W, case in which the optical spectra of the multiwavelength fiber laser consists of 28 equally spaced frequencies covering a range from 192 THz to 197 THz centered in the P23 acetylene absorption line and with a spacing of 200 GHz. As we were increasing the pump power, we could check that the number of frequencies generated did not increase and the optical signal to noise ratio (OSNR) of the frequencies generated decrease.

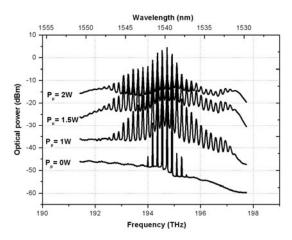


Fig 3.- Output spectra of the laser against different Raman pump power (Raman laser 1455 nm)

In order to increase the number of frequencies generated we introduced another Raman Fiber Laser (RFL2) pumping at $\lambda p_2 = 1427$ nm and output power of Pp₂ = 1W through a 50/50 coupler. Fig 4 shows the new output spectra consisting of 36 equally spaced frequencies covering the frequency range between 192 THz and 198 THz with a spacing of 200 GHz.

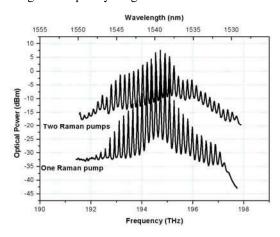


Fig 4.-Comparison of the output spectra of the laser for 1 W pump power for one (1455 nm) and two (1427 & 1455 nm) pump Raman laser

With this configuration we obtain a high resolution and tunable multiwavelength fiber source useful for high precision interrogation of fiber Bragg gratings.

Two mechanisms are responsible of the impossibility to increase the efficiency of the system: the strong backward stimulated Brillouin scattering signal generated in the fiber due to the narrow linewidth of the used lasers (150 kHz) that reduces the FWM generation, and the dispersion properties of the fibers that implies the satisfaction of the phase matching condition only in the zero dispersion wavelength region [11].

3. CONCLUSIONS

In summary, we have experimentally demonstrated a multiwavelength fiber source based on cascaded four-wave mixing in two semiconductor optical amplifiers (SOAs) enhanced by Raman in a dispersion-shifted fiber. The multiwavelength fiber source generated consists of 36 equally spaced and the spacing can be tuned between 50 and 200 GHz (0.4 and 1.6 nm). The source covers a range from 191.56 THz (1565 nm) to 197.88 THz (1515 nm) for a channel spacing of 200 GHz, sweeps with the same spacing over 3 nm and is continuously monitored in frequency along the sweeps. One step to make the source cheaper could be to replace the Tunable Lasers by simple DFB lasers, and act on their temperature and current. This multiwavelength fiber laser source could be useful for simultaneous and high precision interrogation of a network of FBGs.

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